

# Infrared and Raman Spectra of $\text{LiNH}_3\text{OHSO}_4$ and $(\text{NH}_3\text{OH})_2\text{SO}_4$ Single Crystals and Their Deuterated Analogs

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Infrared and Raman spectra of  $\text{LiNH}_3\text{OHSO}_4$  and  $(\text{NH}_3\text{OH})_2\text{SO}_4$  single crystals and their deuterated analogs are recorded and analyzed. Bands are assigned on the basis of  $\text{SO}_4^{2-}$ ,  $\text{NH}_3$ , and OH vibrations. The appearance of the  $\nu_1$  mode in the  $B_g$  species contributing to  $\alpha_{xz}$ ,  $\alpha_{yz}$ , and  $\alpha_{xy}$  polarizability tensor components, the higher wavenumber values of the  $\nu_1$  mode, the observation of the inactive  $\nu_1$  and  $\nu_2$  modes in the infrared spectrum, and the lifting of degeneracies of the  $\nu_2$ ,  $\nu_3$ , and  $\nu_4$  modes suggest that the  $\text{SO}_4$  tetrahedron in  $\text{LiNH}_3\text{OHSO}_4$  is distorted confirming the X-ray data. The large splitting observed for the  $\nu_1$  mode in  $(\text{NH}_3\text{OH})_2\text{SO}_4$  and its decrease on deuteration indicate that the  $\text{SO}_4$  tetrahedron is more distorted in it. Strong hydrogen bonds are formed between the H atoms of the  $\text{NH}_3$  and OH groups with the oxygen atoms of the  $\text{SO}_4$  group in both the crystals. The presence of two bands for the nondegenerate symmetric bending mode ( $\nu_2$ ), the broadness of the stretching modes  $\nu_1$  and  $\nu_3$  of  $\text{NH}_3$ , and the appearance of two  $\text{NH}_3$  torsional bands at 90 K suggest the existence of two distinct  $\text{HONH}_3^+$  sites in the  $(\text{NH}_3\text{OH})_2\text{SO}_4$  crystal confirming the inelastic neutron scattering data. The two  $\text{NH}_3$  torsional modes observed at 90 K suggest that hydrogen bonds formed by  $\text{NH}_3$  protons become more pronounced at low temperature. In  $\text{LiNH}_3\text{OHSO}_4$  only one  $\text{NH}_3$  torsional mode is obtained and this indicates that  $\text{NH}_3\text{OH}^+$  sites in it are identical. ABC bands are observed in the OH stretching region in both of the crystals confirming the presence of strong hydrogen bonds. © 1996 Academic Press, Inc.

## INTRODUCTION

$\text{LiNH}_3\text{OHSO}_4$  crystallizes in the orthorhombic system with space group  $Pbca$  and has eight formula units per unit cell (1). It has a structure which is quite different from the stuffed tetrahedral framework structures of the isoelectronic  $\text{Li}(\text{N}_2\text{H}_5)\text{SO}_4$  and  $\text{Li}(\text{N}_2\text{H}_5)\text{BeF}_4$  or the other compounds of the forms  $\text{LiMSO}_4$  and  $\text{LiMBeF}_4$  ( $M = \text{K}$ ,  $\text{Rb}$ ,  $\text{Cs}$ , and  $\text{NH}_4$ ).  $\text{LiNH}_3\text{OHSO}_4$  contains sheets of  $\text{LiSO}_4$ , hydrogen bonded together by the  $\text{NH}_3\text{OH}^+$  ions. The Li

atom is surrounded by four oxygen atoms at 1.96 Å and S by four oxygen atoms at 1.477 Å.

$(\text{NH}_3\text{OH})_2\text{SO}_4$  crystallizes in the monoclinic system with space group  $P2_1/c$  having four molecular units per unit cell (2). Detailed structure of this compound with bond lengths and bond angles is not available. There are only a few reports of the vibrational analysis of the compounds containing  $\text{NH}_3\text{OH}^+$  ions (3–5). A detailed analysis of the vibrational spectra of  $(\text{NH}_3\text{OH})_2\text{SO}_4$ ,  $(\text{ND}_3\text{OD})_2\text{SO}_4$ ,  $\text{LiNH}_3\text{OHSO}_4$ , and  $\text{LiND}_3\text{ODSO}_4$  has been carried out to obtain more information about the internal structure of these compounds and the nature of hydrogen bonding.

## EXPERIMENTAL

Single crystals of  $\text{LiNH}_3\text{OHSO}_4$  (abbreviated as LNHS) were prepared by the slow evaporation of the aqueous solution containing stoichiometric quantities of  $(\text{NH}_3\text{OH})_2\text{SO}_4$  and  $\text{Li}_2\text{SO}_4$  at room temperature ( $300 \pm 3$  K). Single crystals of  $(\text{NH}_3\text{OH})_2\text{SO}_4$  (abbreviated as NHS) were prepared by the slow evaporation of the aqueous solution of commercially available  $(\text{NH}_3\text{OH})_2\text{SO}_4$  (BDH, AR grade) at room temperature. The deuterated analogs of these compounds were prepared by dissolving them in heavy water (99.99% pure) and evaporating the solution under a vacuum desiccator. The process was repeated several times to enhance the percentage of deuteration.

Well grown single crystals of LNHS and NHS were chosen and crystallographic axes were determined using a polarizing microscope. The sides  $x$ ,  $y$ , and  $z$  of the crystals were chosen in accordance with crystallographic  $a$ ,  $b$ , and  $c$  axes. The sides of the crystals were cut, finely polished, and properly aligned on the goniometer to record the Raman spectra in the  $90^\circ$  scattering geometry. The polarized Raman spectra of LNHS (Figs. 1–3) were recorded on six crystal orientations, viz.,  $y(xx)z$ ,  $x(yy)z$ ,  $x(zz)y$ ,  $y(xy)z$ , and  $y(zy)z$ . Polarized Raman spectra (Figs. 4 and 5) of NHS were recorded on four crystal geometries, viz.,  $z(yy)x$ ,  $z(xz)x$ ,  $z(yz)x$ , and  $z(xy)z$ . Raman spectra of the polycrystalline samples of  $\text{LiNH}_3\text{OHSO}_4$ ,  $\text{LiND}_3\text{ODSO}_4$ ,

<sup>1</sup> To whom correspondence should be addressed.

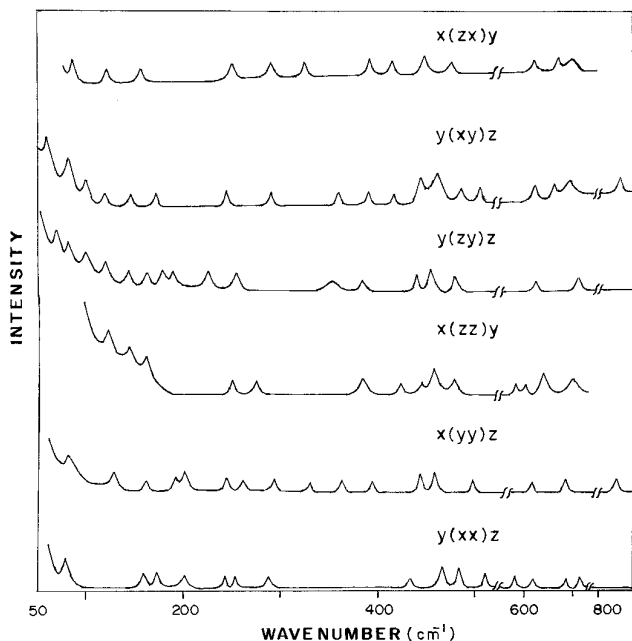


FIG. 1. Raman spectra of  $\text{LiNH}_3\text{OH}\text{SO}_4$  in the 50–850  $\text{cm}^{-1}$  region for  $y(xx)z$ ,  $x(yy)z$ ,  $x(zz)y$ ,  $y(zy)z$ ,  $y(xy)z$ , and  $x(zx)y$  orientations.

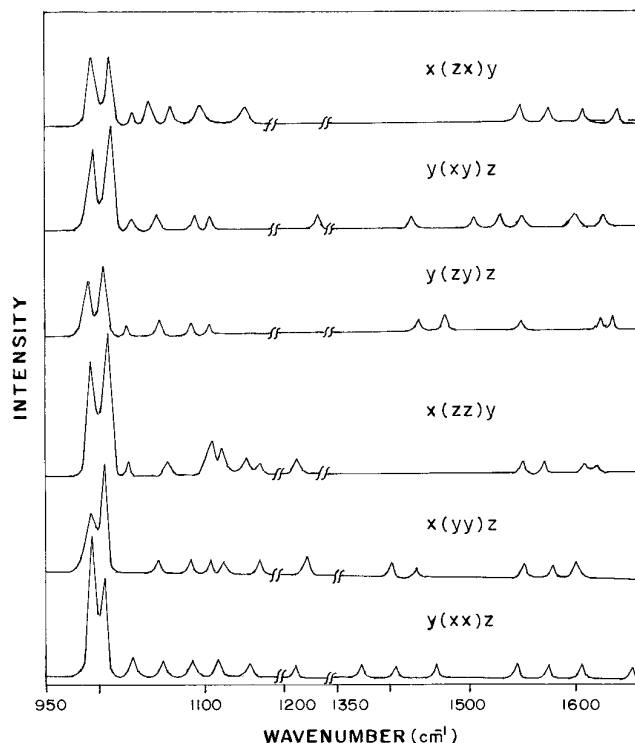


FIG. 2. Raman spectra of  $\text{LiNH}_3\text{OH}\text{SO}_4$  in the 950–1650  $\text{cm}^{-1}$  region for  $y(xx)z$ ,  $x(yy)z$ ,  $x(zz)y$ ,  $y(zy)z$ ,  $y(xy)z$ , and  $x(zx)y$  orientations.

$(\text{NH}_3\text{OH})_2\text{SO}_4$ , and  $(\text{ND}_3\text{OD})_2\text{SO}_4$  were also recorded by taking the samples in capillary tubes. A 1401 Spex Raman spectrometer equipped with a Spectra Physics model 165.08  $\text{Ar}^+$  laser was used for recording Raman spectra. The spectra were recorded using 514.5 nm exciting radiation at a resolution better than 3  $\text{cm}^{-1}$ . The Raman spectra of LNHS and NHS single crystals were recorded in the back scattering geometry,  $x(zz+y)x$  and  $x(yy+z)x$  geometries, respectively, at 300 and 90 K, using a Coherent Innova 300 krypton ion laser (exciting radiation 476.2 nm) and Raman spectra of NHS on a Dilor Z24 Raman spectrometer. The FTIR/IR spectra of LNHS, LNDS, NHS, and a partially deuterated  $(\text{ND}_3\text{OD})_2\text{SO}_4$  (abbreviated NDS) were recorded on a Bruker IFS-66V-FTIR spectrometer/Perkin-Elmer 882 spectrophotometer as polyethylene pellets and KBr pellets. Raman spectra are redrawn conforming to the original spectra obtained for the crystals and are given in Figs. 1–5.

### FACTOR GROUP ANALYSIS

$\text{Li}$ ,  $\text{NH}_3\text{OH}^+$ , and  $\text{SO}_4^{2-}$  ions in LNHS and  $\text{NH}_3\text{OH}^+$  and  $\text{SO}_4^{2-}$  ions in NHS occupy the general site  $C_1$  in the respective crystals. Factor group analysis of the compounds has been carried out using the correlation method developed by Fateley *et al.* (6) (see Tables 1 and 2). The 285 irreducible representations obtained for LNHS excluding acoustic modes at  $k = 0$  are distributed as

$$\Gamma_{\text{LNHS}} = 36A_g + 36B_{1g} + 36B_{2g} + 36B_{3g} + 36A_u + 35B_{1u} + 35B_{2u} + 35B_{3u}.$$

For NHS, 201 optical modes (at  $k = 0$ ) are obtained. They are distributed as

$$\Gamma_{\text{NHS}} = 51A_g + 51B_g + 50A_u + 49B_u.$$

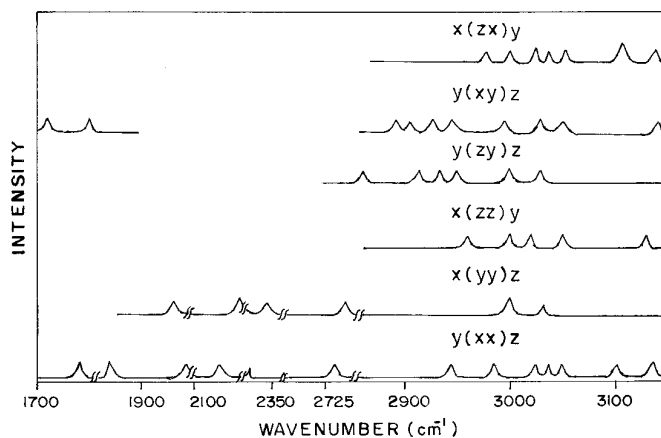


FIG. 3. Raman spectra of  $\text{LiNH}_3\text{OH}\text{SO}_4$  in the 1700–3150  $\text{cm}^{-1}$  region for  $y(xx)z$ ,  $x(yy)z$ ,  $x(zz)y$ ,  $y(zy)z$ ,  $y(xy)z$ , and  $x(zx)y$  orientations.

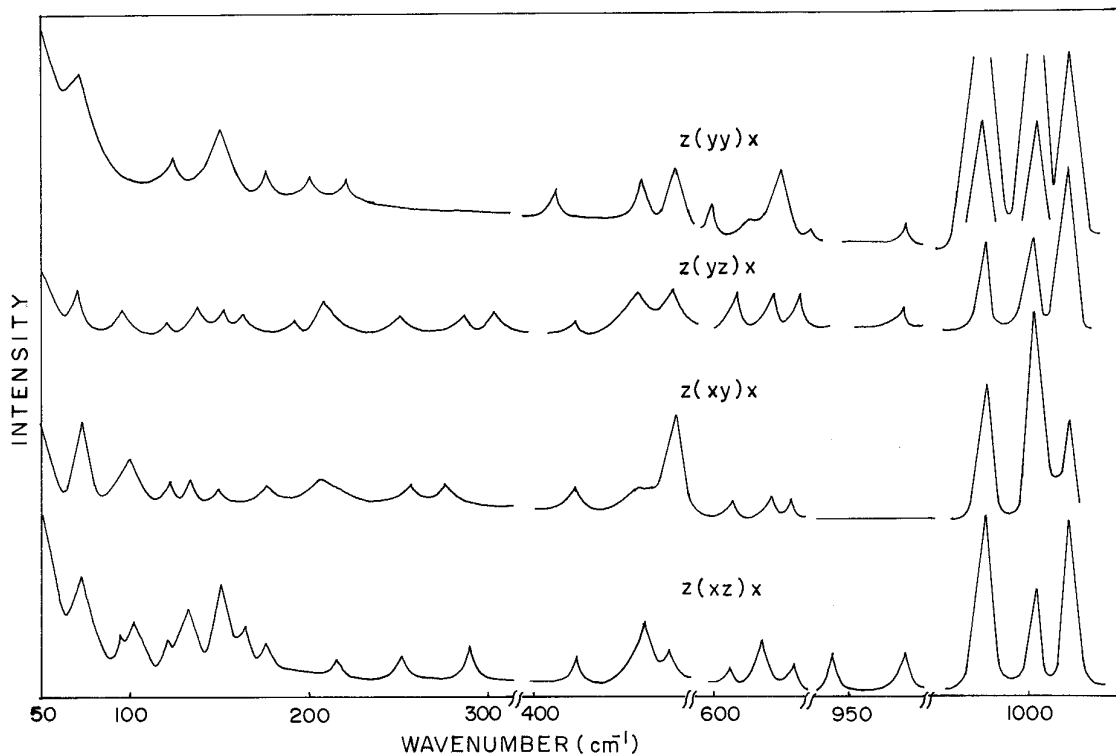


FIG. 4. Raman spectra of  $(\text{NH}_3\text{OH})_2\text{SO}_4$  in the  $50\text{--}1050\text{ cm}^{-1}$  region for  $z(xz)x$ ,  $z(xy)x$ ,  $z(yz)x$ , and  $z(yy)x$  orientations.

## RESULTS AND DISCUSSION

### $\text{SO}_4^{2-}$ Vibrations

A free  $\text{SO}_4^{2-}$  ion under  $T_d$  symmetry has four fundamental vibrations, viz., the nondegenerate symmetric stretching mode  $\nu_1(A_1)$ , the doubly degenerate bending mode  $\nu_2(E)$ , the triply degenerate asymmetric stretching mode  $\nu_3(F_2)$ , and triply degenerate asymmetric bending mode  $\nu_4(F_2)$ . All the modes are Raman active, whereas only  $\nu_3$  and  $\nu_4$  are active in the IR.

### $\text{LiNH}_3\text{OHSO}_4$ and $\text{LiND}_3\text{ODSO}_4$

The nondegenerate stretching mode of  $\text{SO}_4^{2-}$  ( $\nu_1$ ) splits into two in the Raman spectra in all the crystal orientations and in the back scattered spectrum for LNHS (Table 3). In  $y(xx)z$  and  $x(zz)y$  polarizations, two intense bands at  $996$  and  $1011\text{ cm}^{-1}$  are obtained. In all other crystal orientations, these bands have only medium intensity. In  $y(xx)z$  and  $y(xz)x$  orientations, the band at  $996\text{ cm}^{-1}$  shows more intensity than the band at  $1011\text{ cm}^{-1}$ . However, in all other orientations, the band at  $1011\text{ cm}^{-1}$  is more intense than the other band. In the back scattered Raman spectrum, two very intense bands are observed at  $996$  and  $1010\text{ cm}^{-1}$ , with the band at  $1010\text{ cm}^{-1}$  showing enhanced intensity. In the FTIR spectrum, an intense band is observed at

$1001\text{ cm}^{-1}$  for this vibration. In the deuterated compound LNDS, two very intense bands are observed at  $999$  and  $1009\text{ cm}^{-1}$  in the Raman spectrum and two weak bands at  $980$  and  $1000\text{ cm}^{-1}$  in the IR spectrum.

The asymmetric stretching mode  $\nu_3$  of  $\text{SO}_4^{2-}$  is observed with degeneracy lifted in the Raman spectra in all the crystal orientations and in the back scattered spectrum (Table 3). In the FTIR spectrum, a very intense band is observed at  $1132\text{ cm}^{-1}$  for this mode. In the deuterated compound, degeneracy is lifted in both the IR and Raman spectra.

The  $\nu_2$  mode of  $\text{SO}_4^{2-}$  appears with degeneracy lifted in the Raman spectra in all the crystal orientations. However, in the back scattered Raman spectrum and in the FTIR spectrum only a single band is observed. The asymmetric bending mode  $\nu_4$  also appears with degeneracy lifted in the Raman spectra, whereas it is partially retained in the FTIR spectrum.

Under  $T_d$  symmetry, the  $\nu_1$  mode has polarizability components  $\alpha_{xx}$ ,  $\alpha_{yy}$ , and  $\alpha_{zz}$  and it is expected to appear in the  $A_g$  orientations without any distortion of the  $\text{SO}_4$  tetrahedra. This mode is also observed in  $B_{1g}$ ,  $B_{2g}$ , and  $B_{3g}$  orientations with considerable intensity. Since the  $\text{SO}_4^{2-}$  ion occupies a site  $C_1$ , it can be seen from the correlation table that this mode can be active in the  $B_{1g}$ ,  $B_{2g}$ , and  $B_{3g}$  species, due to the distortion of  $\text{SO}_4$  tetrahedra from  $T_d$

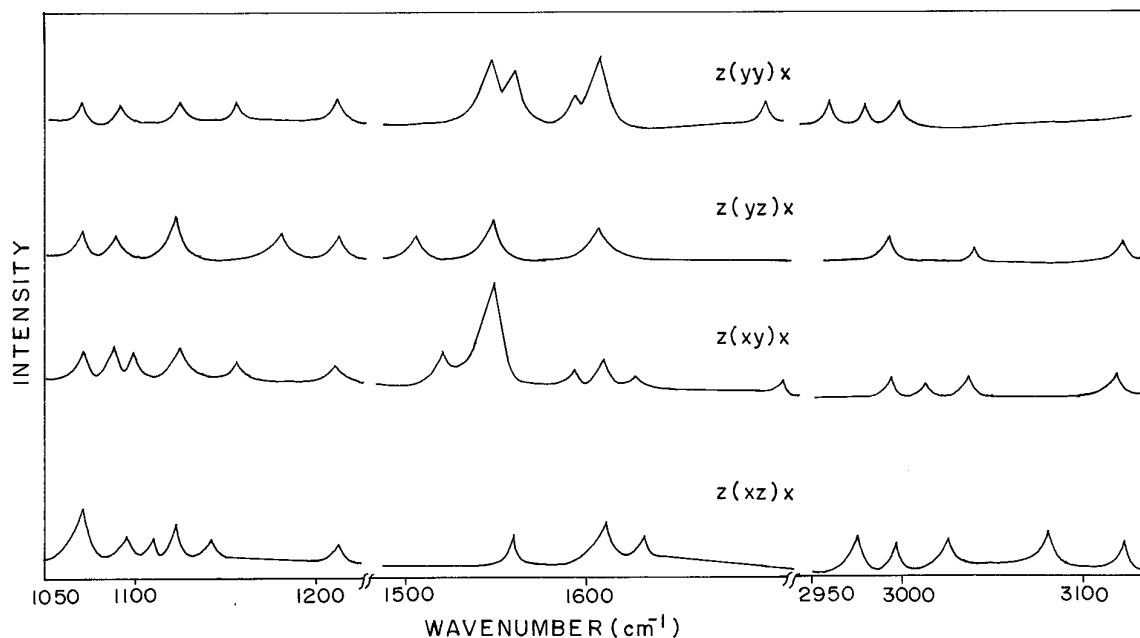


FIG. 5. Raman spectra of  $(\text{NH}_3\text{OH})_2\text{SO}_4$  in the 1050–3150  $\text{cm}^{-1}$  region for  $z(xz)x$ ,  $z(xy)x$ ,  $z(yz)x$ , and  $z(yy)x$  orientations.

to  $C_1$ , contributing to  $\alpha_{xz}$ ,  $\alpha_{yz}$ , and  $\alpha_{xy}$  polarizability tensor components (7). Higher wavenumber values obtained for the nondegenerate stretching mode  $\nu_1$  than those in a free  $\text{SO}_4^{2-}$  ion (8) also confirms the distortion of  $\text{SO}_4$  tetrahedra as is evident from different S–O bond lengths determined from the X-ray diffraction studies (1, 9). Distortion of the  $\text{SO}_4^{2-}$  ion and the fact that there are eight molecular units in the Bravais cell leads to a splitting of the  $\nu_1$  mode and additional splitting, apart from the lifting of degeneracies of the  $\nu_2$ ,  $\nu_3$ , and  $\nu_4$  modes. X-ray data show that H atoms of the  $\text{NH}_3$  group and OH group form strong hydrogen bonds with the oxygen atoms of the  $\text{SO}_4$  group with the N–H  $\cdots$  O distance ranging from 2.799 to 3.066 Å and the O–H  $\cdots$  O distance 2.65 Å. The presence of these hydrogen bonds may be the reason for the observed distortion in the  $\text{SO}_4$  tetrahedra in LNHS. Deuteration does not affect the  $\text{SO}_4^{2-}$  vibrations.

#### $(\text{NH}_3\text{OH})_2\text{SO}_4$ and $(\text{ND}_3\text{OD})_2\text{SO}_4$

The nondegenerate symmetric stretching mode of the  $\text{SO}_4^{2-}$  ion exhibits larger splitting in both the FTIR and Raman spectra of NHS. Three very intense bands around 981, 1005, and 1021  $\text{cm}^{-1}$  are obtained for this mode in  $z(yy)x$  polarization of the Raman spectra of the crystal and in the powder spectrum (Table 4). Intensity of these bands alternates in the other polarizations. In the FTIR spectrum, only two intense bands at 978 and 1012  $\text{cm}^{-1}$  are obtained. In the deuterated sample, three bands are again obtained but at 984, 992, and 1002  $\text{cm}^{-1}$  in the Raman spectrum. In the FTIR spectrum of NDS, two intense bands

are observed at 977 and 995  $\text{cm}^{-1}$  for this mode. In NHS, the  $\nu_1$  mode has a frequency spread of 40  $\text{cm}^{-1}$  in Raman and 34  $\text{cm}^{-1}$  in the IR, while in NDS, this is only 18  $\text{cm}^{-1}$  in both the spectra. This large splitting observed for the  $\nu_1$  mode in NHS and its decrease in NDS indicate that the  $\text{SO}_4$  tetrahedron is distorted in NHS. Hydrogen atoms of  $\text{NH}_3$  and OH groups form hydrogen bonds in the NHS crystal. Deuteration of the compound causes a reduction of the hydrogen bond strength, leading to a decrease in the spread of the splitting of this mode.

The degeneracies of the  $\nu_2$ ,  $\nu_3$ , and  $\nu_4$  modes are lifted in the Raman spectra of the NHS crystal and of the deuterated compounds (Table 4). This can be expected from the lowering of the symmetry of the  $\text{SO}_4^{2-}$  ion from  $T_d$  to  $C_1$ . The appearance of the  $\nu_1$  mode in the  $B_g$  species also confirms the distortion of the  $\text{SO}_4$  tetrahedra (7). The observed large splitting of the  $\nu_1$  mode in NHS compared to LNHS suggests that the  $\text{SO}_4^{2-}$  ion is more distorted in NHS.

#### $\text{NH}_3\text{OH}^+$ Vibrations

$\text{NH}_3\text{OH}^+$  vibrations consist of  $\text{NH}_3$  and OH vibrations.  $\text{NH}_3$  in the free state has  $C_{3v}$  symmetry and it has four fundamental vibrations, all being both Raman and IR active. The symmetric [ $\nu_1(A_1)$ ] and asymmetric [ $\nu_3(E)$ ] N–H stretching modes appear with weak intensity in the Raman spectra of all the compounds. In the FTIR spectrum of NHS, a very intense broad band is observed around 2955  $\text{cm}^{-1}$  for both the stretching modes. Deuteration of this compound is only partial (about 70%) as seen from the stretching bands of  $\text{ND}_3$  in NDS. These modes shift to

TABLE 1  
Correlation of the Internal Vibrational Modes of  
 $\text{SO}_4^{2-}$  and  $\text{NH}_3$  in  $\text{LiNH}_3\text{OH}\text{SO}_4$

$\text{SO}_4^{2-}$		
Free ion symmetry	Site symmetry	Factor group symmetry
$T_d$	$C_1$	$D_{2h}$
8 $A_1$	A	$A_g$ 9
		$A_u$ 9
16 $E$		$B_{1g}$ 9
		$B_{2g}$ 9
		$B_{3g}$ 9
		$B_{1u}$ 9
		$B_{2u}$ 9
		$B_{3u}$ 9
48 $F_2$		
$\text{NH}_3$		
Molecular symmetry	Site symmetry	Factor group symmetry
$C_{3v}$	$C_1$	$D_{2h}$
16 $A_1$	A	$A_g$ 6
		$A_u$ 6
		$B_{1g}$ 6
		$B_{2g}$ 6
		$B_{3g}$ 6
		$B_{1u}$ 6
		$B_{2u}$ 6
		$B_{3u}$ 6
32 $E$		

2267  $\text{cm}^{-1}$  with shoulders at 2150 and 2375  $\text{cm}^{-1}$  in the FTIR spectrum of the deuterated compound. In the Raman spectrum, they exhibit weak bands at 2155 and 2260  $\text{cm}^{-1}$ . In the FTIR spectrum of LNHS, these modes obtained as an intense broad band around 2930  $\text{cm}^{-1}$  are observed as two bands at 2285 and 2318  $\text{cm}^{-1}$ .

The Raman bands appearing in the 1500–1650  $\text{cm}^{-1}$  region in both NHS and LNHS are assigned to the bending modes  $\nu_2$  and  $\nu_4$  of  $\text{NH}_3$ . One intense band each is observed for  $\nu_2$  and  $\nu_4$  modes in the IR spectra for both LNHS and its deuterated analog. But in NHS, the nondegenerate  $\nu_2$  ( $\text{NH}_3$ ) mode splits into two bands of equal intensity in both FTIR and Raman spectra while the degeneracy of the  $\nu_4$  mode is lifted in all the Raman orientations and in the FTIR spectrum.

From a study of the inelastic neutron scattering spectra, Jayasooriya *et al.* (5) predicted the existence of at least two distinct  $\text{HONH}_3^+$  sites in the NHS crystal from two  $\text{NH}_3$  torsional vibrations about the N–O axis. They have also obtained two IR  $\text{NH}_3$  torsional modes at liquid nitrogen temperature. In the present study, the band at 416  $\text{cm}^{-1}$  observed in the FTIR spectrum at room temperature is assigned to this mode as it disappears on deuteration. Weak bands are also observed in this region in the  $z(yy)x$ ,  $z(xz)x$ , and  $z(xy)x$  polarizations of the Raman spectra. Even though only one band is observed at room tempera-

ture for the torsional mode, the appearance of two bands in  $\nu_2\text{NH}_3$  mode region and the broadness of the  $\nu_1$  and  $\nu_3$  stretching modes can be considered as a confirmation of the existence of two distinct  $\text{HONH}_3^+$  sites in the NHS crystal. Two bands are also obtained in NDS for the  $\nu_2$  mode with a higher wavenumber band at 1427  $\text{cm}^{-1}$  having a higher intensity than the other (1387  $\text{cm}^{-1}$ ).

The N–H stretching frequencies  $\nu_1$  and  $\nu_3$  appear at lower wavenumbers than those in a free  $\text{NH}_3$  group (8, 10) indicating that strong hydrogen bonds are formed by the H atom of the  $\text{NH}_3$  group with the oxygen atoms of the  $\text{SO}_4^{2-}$  anion in both LNHS and NHS crystals. This is consistent with the X-ray structural data of N–H  $\cdots$  O distances of 2.799–3.066 Å obtained for LNHS (1).

Two weak bands are observed at 2000 and 2059  $\text{cm}^{-1}$  in the FTIR spectra of NHS. They are assigned to combination bands of the  $\text{NH}_3$  deformations (1614  $\text{cm}^{-1}$ ) and  $\text{NH}_3$  torsional modes (380 and 430  $\text{cm}^{-1}$ ). Similar combination bands have been observed earlier in  $\text{N}_2\text{H}_6\text{SO}_4$  crystals (11). In LNHS, only one band around 1950  $\text{cm}^{-1}$  is obtained in this region.

The positive isotope effect implies a weakening of the hydrogen bonds on deuteration (12, 13). The isotopic frequency ratios observed for the N–H stretching modes for NHS and LNHS are about 1.289 and 1.267, respectively. For the bending modes, these ratios have values of 1.10156 and 1.1868, respectively. For a free molecule, this ratio is about 1.355. This value decreases as the strength of the hydrogen bond increases. Therefore, the lower values obtained for this ratio also indicate the presence of strong hydrogen bonds in  $\text{LiNH}_3\text{OH}\text{SO}_4$  and  $(\text{NH}_3\text{OH})_2\text{SO}_4$ .

TABLE 2  
Correlation of the Internal Vibrational Modes of  
 $\text{SO}_4^{2-}$  and  $\text{NH}_3$  in  $(\text{NH}_3\text{OH})_2\text{SO}_4$

$\text{SO}_4^{2-}$		
Free ion symmetry	Site symmetry	Factor group symmetry
$T_d$	$C_1$	$C_{2h}$
4 $A_1$	A	$A_g$ 9
		$B_g$ 9
8 $E$		$A_u$ 9
		$B_u$ 9
24 $F_2$		
$\text{NH}_3$		
Molecular symmetry	Site symmetry	Factor group symmetry
$C_{3v}$	$C_1$	$C_{2h}$
16 $A_1$	A	$A_g$ 12
		$B_g$ 12
		$A_u$ 12
32 $E$		$B_u$ 12

TABLE 3  
Spectral Data ( $\text{cm}^{-1}$ ) and Band Assignments for  $\text{LiNH}_3\text{OH}\text{SO}_4$  and  $\text{LiND}_3\text{ODSO}_4$

$\text{LiNH}_3\text{OH}\text{SO}_4$													
IR	FTIR	Raman						Back scattered		$\text{LiND}_3\text{ODSO}_4$		Assignments	
		$y(xx)z$	$x(yy)z$	$x(zz)y$	$x(zx)y$	$y(xy)z$	$y(zy)z$	300 K	90 K	Raman	IR		
		3	4	5	6	7	8	9	10	11	12		13
	63vw 87vs 102w 131m 141vw		77w 79w 131w	120w	121w	101vw 101vw 121w 146vw	56w 70w 81m 81w 101vw 101vw 121w 141vw		116w 115w	94w 148w 167vw 207vw 226vw	80w 130w 140w 174vw 221w	$\text{SO}_4$ translation $\delta(\text{O} \cdots \text{O})$ $\text{SO}_4$ libration $\text{NH}_3\text{OH}$ translation	
	150w 169m 198s 240sh 255vw 286sbr 333vs	159w 170vw 201w 246w 255vw 285w	165vw 191w 201w 246w 259w 290w 330vw 365w	161vw 161vw 251w 250w 266w	161vw 166vw 285vw 325w	166vw 161vw 241vw 226vw 261w		230wbr 245vwbr 290w	226vw 248vw	221w 240w	240w	$\nu(\text{O} \cdots \text{O})$ $\text{Li}^+$ translation Li-O stretch $\text{NH}_3$ torsion	
390w	388vs 422vs		390w 434vw 425vw	386w 425vw	390w 415vw	391vw 415vw	361vw 360vwbr 381vw			365w 430m	380w 408w	408w	Li-N stretch Li-O stretch $\text{NH}_3$ torsion
440w 472m	422vs 474vs	463wbr 481w	440w 455w	445w 454m 481w	456w 481w	449w 456w 486vw	441vw 451w 481w	450m	452m 486m	444m 450m 461vw 476w	444m 482w	468m 482w	$\nu_2\text{SO}_4$
520vw 621s	600sh 639vs	505w 585vw 615vw 643w 651w	491w 591vw 615vw 641vw	591vw 601vw 621m 651wbr	615vw 638w 635vw 651w	615vw 615vw 635vw 655w		618mbr 635w 648m	503vw 617m 638w 647wbr	610sh 625w 635m 651w	519w 618s 642sh	618s	Li-N stretch $\nu_2\text{SO}_4$
752w	800sh		820w										$\gamma(\text{OH})$ $\text{NH}_3/\text{ND}_3$ rock $2\nu_2\text{SO}_4$
995m	1001s	996s 1011s 1030vw	996m 1010s 1030w	996s 1011s 1030w	996m 1011m 1030w	998m 1010m 1030w	997m 1011m 1031vw	996vs 1010vs 1038vw	997s 1011vvs	996vs 774vw	980w 1000w	980w	$\nu_1\text{SO}_4$ N-OH/N-OD stretch
1083w 1087vs 1121vs	1132vsbr	1061w 1086w 1111w 1145vw 1156w	1061w 1086vw 1103m 1106w 1113w 1140w 1143w	1066w 1066w 1103m 1101w 1101w 1143w	1066w 1066w 1103m 1101w 1101w	1068w 1068w 1091w 1091vw 1106vw	1083sh 1087m 1102s 1102s 1138w 1143m 1153w 1157m 1177w	1087m 1102s 1102s 1138w 1143m 1153w 1157m 1177w	1052vw 1087m 1102s 1143m 1153w 1157m 1177w	1052vw 1092vs 1117vw 1141vw 1151vw	1050m 1092vs 1127vs 1162s	1050m	$\nu_3\text{SO}_4$
1245mbr		1210vw	1220w	1210w		1206vw		1214w	1218mbr	922vw 960vw	923m 1032m	923m	$\delta\text{OH}/\text{OD}$
1405vw		1375w 1405w 1440w	1430vw			1425vw	1430vw 1455w						Combinations
1525s 1565s	1506s	1540w 1561w 1580w	1556vw 1581w	1545w 1566w	1551w 1571w	1516vw 1526vw 1548w	1550w	1516w 1535vw	1521vw 1534vw	1193m 1275w	1232m	1232m	$\nu_2\text{NH}_3/\text{ND}_3$
1640s	1614s	1605vw 1660w	1600w	1616w	1611w 1640w	1600w 1626w 1676vw	1625vw 1635vw	1618mbr	1615w 1624m	1376w 1430w	1360w 1422w	1360w	$\nu_4\text{NH}_3/\text{ND}_3$
	1760w	1746w 1870w 2125vw	2140w			1711w 1760vw				1615vw 1660vw	2165w	2165w	Combinations
1935wbr	1950w	1945w	1930w										$\text{NH}_3$ torsion- $\text{NH}_3$ deformation
2230vw 2725m	2365w 2697vs	2330w 2730wbr	2340vw 2740w							2804w 2832m 2859m	2060w	1750w 2060w	$\nu(\text{OH})$ A & B bands Combinations
2920m 3030s	2930sbr	2945w 2985w 3020w 3030w 3035w 3050vw 3100vw	2995w 3030w	2960w 3020w 3020w 3030w 3052w 3130w	2970w 3020w 3030w 3030w 3050w 3110w 3140w	2890w 2905w 2930vw 2945w 2950vw 3000vw 3030w	2860w 2915vw 2935vw 2950vw 3000vw 3030w	2916w 2964w 3010w 3050w	2872w 2900w 2935s 2954w 3012w 3086w 3122w 3152w 3259w	2176w 2260s 2360sbr 2990vw 3014vw	2285s 2318m 2390m 2941vw 3055vw	2390m	$\nu_1$ & $\nu_3$ $\text{NH}_3/\text{ND}_3$
3140vw													Combinations

Note. vs, very strong; s, strong; m, medium; vw, very weak; w, weak; br, broad; sh, shoulder.

TABLE 4  
Spectral Data (cm<sup>-1</sup>) and Band Assignments

(NH <sub>3</sub> OH) <sub>2</sub> SO <sub>4</sub>									
Raman				Back scattered		(ND <sub>3</sub> OD) <sub>2</sub> SO <sub>4</sub>			
<i>z</i> ( <i>yy</i> ) <i>x</i>	<i>z</i> ( <i>xz</i> ) <i>x</i>	<i>z</i> ( <i>yz</i> ) <i>x</i>	<i>z</i> ( <i>xy</i> ) <i>x</i>	300 K	90 K	FTIR	FTIR	Raman	Assignments
1	2	3	4	5	6	7	8	9	10
72w	72m 92w 100w	72w 98w	74m 96mbr	80w 95vw 107vw	69m 84w 96w	87vs	60vw 69vw 87w	85w	SO <sub>4</sub> translation
124w 150m	120vw 130m	120vw 135w	120vw 130vw	139w 158m	111vw 135m 160w	125s 150mbr	125sh 150w	126m 150w	SO <sub>4</sub> libration $\delta(\text{O} \cdots \text{O})?$
175vw	165sh 175vw	160w 192w	170vw	175m	177m 192vw	190vsbr	190sh	195w	SO <sub>4</sub> libration
200vwbr 215w	215vw 250w	205mbr 250w	200wbr 255vw	212mbr	204m 217m 223vw 259vw	230vs	218vs 265sh		NH <sub>3</sub> OH translation
	285w	286w 300w	270vw	259w 288w	289m 306w 380m	290sh		274w 295vw	$\nu(\text{O} \cdots \text{O})$
410w	420w		420w		424m	417vs			NH <sub>3</sub> torsion
460w 475m	458m 475w	455m 475m	455sh 478s	454m 475m	460m 473m	459m 471m	468s	456m 478m	$\nu_2\text{SO}_4$ OD torsion?
595w 620sh 637m	605w 625m 640vw	610w 630w 641w	610vw 630w 638vw	603w 626m 638sh	601w 619m 627m 637m 656vw	603vs 615vs 627vs	577s 610vs 635vs	620w 643w	$\nu_4\text{SO}_4$
655vw									OH torsion $\gamma\text{OH}/\text{OD}$
						831s	650sh 835w	665vw	
880vw	835w 880w	880vw		842w	850m				NH <sub>3</sub> rocking
981vs 1005vs 1021vs	981s 1003m 1021s	983m 1004m 1020s	981s 1004vs 1021s	979vs 1002s 1017vs	978vs 999s 1017vs	978s 1012s	977s 995s	984vs 992vs	$\nu_1\text{SO}_4$
1070w 1090w 1122vw 1160vw	1070m 1090w 1110vw 1120w 1150w	1070w 1089vw 1122m 1185w	1070vw 1085w 1094w 1120w 1160vw	1067w 1084w 1115w	1067m 1086m 1102w 1121m	1066vs 1120sh	1097vsbr 1183sh	1075wbr 1110w 1170w	$\nu_3\text{SO}_4$
1215w	1215w	1215w	1210wbr	1208wbr	1217mbr	1215s	892m 931m	950w	$\delta\text{N-OH}/\text{N-OD}$
1550m 1560w	1560vw 1550w	1505w 1550w	1520w 1550m	1506vw 1550w	1508w 1552w	1523s 1558s	1387m 1427m 1529m 1557m	1265vw 1330m	$\nu_2\text{NH}_3/\text{ND}_3$
1594w 1608m 1700vw	1610w 1635vw	1610wbr	1590vw 1610w 1625vw 1713vw	1608w	1583vw 1614m	1616s 1654sh	1467m 1616w	1350w 1390w	$\nu_4\text{NH}_3/\text{ND}_3$
						2000w 2059w	1874vw		NH <sub>3</sub> deformation + NH <sub>3</sub> torsion

TABLE 4—Continued

$(\text{NH}_3\text{OH})_2\text{SO}_4$									
Raman				Back scattered		FTIR	$(\text{ND}_3\text{OD})_2\text{SO}_4$		Assignments
$z(yy)x$	$z(xz)x$	$z(yz)x$	$z(xy)x$	300 K	90 K		FTIR	Raman	
1	2	3	4	5	6	7	8	9	10
						2100w 2230w 1684sh 2400sh			Combinations
				2737wbr	2710vw 2750vw	2739vs	1765w 2055w 2793m		$\nu(\text{OH})$ ABC bands
2960w 2980w	2975w 2990w	2990w	2990w	2963m	2911vw 2928w	2955vs	2150sh 2887s	2115w 2145w 2155w	$\nu_1\text{NH}_3/\text{ND}_3$
2999w	3030w 3080w 3120vw	3120vw	3015vw 3040vw 3120vw	3011m 3122m 3200w	3019w 3032w 3063vw 3129w		2267vs 2375sh	2260w	$\nu_2\text{NH}_3/\text{ND}_3$

Three fundamental vibrations connected with the hydrogen bonded OH groups are  $\nu(\text{OH})$  stretching and inplane  $\delta(\text{OH})$  and out-of-plane  $\gamma(\text{OH})$  bending vibrations. In a strongly hydrogen bonded system, the  $\nu\text{OH}$  band is usually very broad due to strong interactions between the proton vibrations and the  $\nu(\text{O} \cdots \text{O})$  vibrations (14, 15). Due to Fermi resonance, the overtones of the  $\delta\text{OH}$  and  $\gamma\text{OH}$  modes interact with the broad  $\nu\text{OH}$  band leading to the appearance of three bands A, B, and C in the  $\nu\text{OH}$  stretching region (16). In  $\text{LiNH}_3\text{OH}\text{SO}_4$ , strong hydrogen bonds are present between H atoms of the OH group of the  $\text{NH}_3\text{OH}^+$  ion and O atoms of the  $\text{SO}_4$  groups with  $\text{O} \cdots \text{O}$  distances of around 2.65 Å. In  $\text{NH}_4\text{HSeO}_4$ , A, B, and C bands are obtained at 2655, 2290, and 1610  $\text{cm}^{-1}$ , respectively, where the  $(\text{O} \cdots \text{O})$  distance is around 2.56 Å (17, 18). In  $\text{CsHSeO}_4$ , with  $(\text{O} \cdots \text{O})$  distances around 2.605 Å, they are observed at 2750, 2380, and 1600  $\text{cm}^{-1}$ , respectively (19). In the IR and Raman spectra of the powder (20) samples of  $\text{LiNH}_3\text{OH}\text{SO}_4$  and  $\text{LiND}_3\text{ODSO}_4$  reported earlier, the bands at 3110, 2725, and 1935  $\text{cm}^{-1}$  have been assigned to the ABC triplets, but in the FTIR spectrum no band is obtained around 3110  $\text{cm}^{-1}$ . Therefore A, B, and C bands in the FTIR and Raman spectra are reassigned by considering the  $\text{O} \cdots \text{O}$  distance (2.65 Å) in this crystal.

In the FTIR spectrum of LNHS, a broad intense band from 3150 to 2850  $\text{cm}^{-1}$  with a peak at 2930  $\text{cm}^{-1}$  is obtained. This broad band must be due to the symmetric and asymmetric N–H stretching modes as the H atoms in  $\text{NH}_3$  group also form hydrogen bonds. Therefore the A and B

bands are assigned at 2694 and 2363  $\text{cm}^{-1}$ , respectively. The C band cannot be identified as it is masked by the  $\text{NH}_3$  bending modes. In LNDS, the A and B bands shift to 2060 and 1750  $\text{cm}^{-1}$  in the IR spectrum.

In  $(\text{NH}_3\text{OH})_2\text{SO}_4$ , the A and B bands are not well resolved. A broad band extending from 2800 to 2200  $\text{cm}^{-1}$  appears in the FTIR spectrum, with a peak at 2739  $\text{cm}^{-1}$  and a shoulder at 2400  $\text{cm}^{-1}$ . The C band is also observed as a shoulder at 1684  $\text{cm}^{-1}$ . In the Raman spectrum, the A band is observed as a weak one at 2737  $\text{cm}^{-1}$ . In the deuterated compound  $(\text{ND}_3\text{OD})_2\text{SO}_4$ , the A and B bands are observed at 2055 and 1765  $\text{cm}^{-1}$ , respectively. However, the C band is not identified as it is masked by the strong S–O stretching modes. The  $\delta(\text{OH})$  and  $\gamma(\text{OH})$  modes are also assigned in both the LNHS and NHS (Tables 3 and 4).

The isotopic frequency ratio for A and B bands are of the order obtained for the  $\text{NH}_3$  stretching modes indicating strong hydrogen bonds of the type  $\text{O}–\text{H} \cdots \text{O}$  in the crystal. The  $(\text{O} \cdots \text{O})$  stretching mode is observed around 290  $\text{cm}^{-1}$  in the Raman spectra of LNHS. On deuteration this band shifts to 240  $\text{cm}^{-1}$ . This mode is obtained as a weak band around 288  $\text{cm}^{-1}$  in the Raman spectra and as a shoulder around 290  $\text{cm}^{-1}$  in the FTIR spectrum of NHS. A medium intense band is observed around 150  $\text{cm}^{-1}$  in most of the crystal orientations of the Raman spectra of NHS. In the FTIR spectrum, a broad band is also obtained in this region. On deuteration the intensity of this band is considerably reduced. Hence, this band is assigned to the  $\delta(\text{O} \cdots \text{O})$  mode (15).

The –OH torsional modes occur at considerably higher



frequencies, owing to the shorter O–H bond and much lower moment of inertia compared to the NH<sub>3</sub> group. In the inelastic neutron scattering study, –OH torsional mode appears as a broad band around 656 cm<sup>-1</sup> (5). In the FTIR spectrum of NHS, this mode is not identified as they are masked by the strong asymmetric bending vibration of SO<sub>4</sub><sup>2-</sup>. However, in the deuterated compound a new mode is observed at 577 cm<sup>-1</sup> and is assigned to the –OD torsional modes.

### External Modes (Lattice Modes)

Lattice modes are observed below 400 cm<sup>-1</sup>. An unambiguous assignment of bands in this region is difficult as a large number of bands due to translational and rotational modes of anions and cations appears there. The translational mode of the Li<sup>+</sup> ion occurs at a higher wavenumber than those of other lattice modes in the 300–400 cm<sup>-1</sup> region (21). These bands will not be observed in NHS. Therefore the bands observed around 365 cm<sup>-1</sup> in a few crystal orientations of the Raman spectra are assigned to the translational modes of the Li<sup>+</sup> cation (22). The bands around 380 cm<sup>-1</sup> are due to Li–O stretching modes (23, 24). These bands are not affected by deuteration. The Raman band around 87 cm<sup>-1</sup> in NHS remains intact on deuteration. This suggests that these bands are due to the translational modes of SO<sub>4</sub><sup>2-</sup> (25).

### Temperature Effects

Back scattered Raman spectra of single crystals of LiNH<sub>3</sub>OHSO<sub>4</sub> and (NH<sub>3</sub>OH)<sub>2</sub>SO<sub>4</sub> recorded at 300 and 90 K show significant variations.

*LiNH<sub>3</sub>OHSO<sub>4</sub>.* Two very intense bands are observed at 300 K around 995 and 1010 cm<sup>-1</sup> in the  $\nu_1$  mode region of the SO<sub>4</sub><sup>2-</sup> ion. At 90 K, the band at 995 cm<sup>-1</sup> shifts to 997 cm<sup>-1</sup> but with reduced intensity while the other band around 1011 cm<sup>-1</sup> appears with an enhancement of the intensity. Both these bands become sharp at low temperature. In the asymmetric SO<sub>4</sub><sup>2-</sup> stretching region the shoulder at 1083 cm<sup>-1</sup> disappears and a medium intense band appears around 1087 cm<sup>-1</sup> at 90 K. The intensity of the band at 1102 cm<sup>-1</sup> is enhanced and the weak bands at 1138 and 1153 cm<sup>-1</sup> are shifted to 1143 and 1157 cm<sup>-1</sup> with enhanced intensity. A new band is also observed at 1177 cm<sup>-1</sup>.

The single band at 450 cm<sup>-1</sup> appearing for the  $\nu_2$  mode in the back scattered spectrum recorded at 300 K splits into two bands at 452 and 485 cm<sup>-1</sup> in the low temperature spectra. In the  $\nu_4$  mode, intensity sharing of the bands is observed. In the room temperature spectra, the band at 617 cm<sup>-1</sup> is broader and less intense than the bands at 648 cm<sup>-1</sup>. At 90 K, the band at 618 cm<sup>-1</sup> becomes sharper and gains intensity, while the band at 647 cm<sup>-1</sup> becomes less intense but broad.

In the NH<sub>3</sub> stretching mode region bands at 90 K are

shifted to lower wavenumbers indicating the strengthening of hydrogen bonds. Splitting of bands and considerable increase in intensity of bands are also observed. These indicate Fermi resonance between NH<sub>3</sub> stretching modes and the combinations and overtones of NH<sub>3</sub> deformation vibrations (26). This implies that hydrogen bonds become stronger at low temperature.

Torsional modes of NH<sub>3</sub>, which were absent at the room temperature back scattered Raman spectra, appear at 430 cm<sup>-1</sup> at 90 K. This band is assigned in comparison with the corresponding bands observed in NHS. As explained in the following section of NHS, the appearance of this band indicates the strengthening of hydrogen bonds in LNHS at 90 K.

*(NH<sub>3</sub>OH)<sub>2</sub>SO<sub>4</sub>.* In the spectra recorded at room temperature, three bands are obtained for the  $\nu_1$  mode of SO<sub>4</sub><sup>2-</sup>. The band at 979 cm<sup>-1</sup> has a higher intensity than the band at 1017 cm<sup>-1</sup>. At 90 K, the Raman band at 978 cm<sup>-1</sup> has lesser intensity than the band at 1017 cm<sup>-1</sup>. Thus intensity sharing takes place among the bands in NHS. In the  $\nu_3$  mode region, bands observed at 90 K are at higher wavenumbers and at larger intensity than those observed at 300 K. A new band is also observed at 1102 cm<sup>-1</sup>. The sharing of intensity between bands for the  $\nu_2$  mode is also obtained in NHS. Moreover, the separation between the  $\nu_2$  bands reduces to 13 cm<sup>-1</sup> at 90 K from 21 cm<sup>-1</sup> at 300 K. In the asymmetric bending mode ( $\nu_4$ ) region the medium intense band at 626 cm<sup>-1</sup> splits into two medium intense bands at 619 and 627 cm<sup>-1</sup> and the shoulder at 638 cm<sup>-1</sup> becomes well resolved to a band at 637 cm<sup>-1</sup>.

As explained earlier, Jayasooriya *et al.* (5) observed two NH<sub>3</sub> modes at 380 and 425 cm<sup>-1</sup> in the IR spectrum recorded at liquid nitrogen temperature. In the present study two medium intense bands are obtained at 380 and 424 cm<sup>-1</sup> in the back scattered Raman spectrum at 90 K, confirming the observation of Jayasooriya *et al.* (5) of the existence of two distinct NH<sub>3</sub>OH<sup>+</sup> sites in the crystal lattice. It may be noted that these bands are not observed in the back scattered Raman spectra at room temperature even though one band each is observed in certain polarizations of the Raman spectra. At low temperature, hydrogen bonds formed by the NH<sub>3</sub> protons become more pronounced and the distinct nature of two NH<sub>3</sub>OH<sup>+</sup> sites becomes more clear leading to the appearance of two NH<sub>3</sub> torsional modes. This also confirms that the NH<sub>3</sub>OH<sup>+</sup> sites in LNHS are identical since only one band is obtained for the NH<sub>3</sub> torsional mode.

Bands corresponding to the  $\nu(\text{O} \cdots \text{O})$  and  $\delta(\text{O} \cdots \text{O})$  modes appear at 90 K with enhanced intensity indicating the strengthening of hydrogen bonds.

In the Raman spectrum recorded at 90 K, the OH torsional mode is observed as a very weak broad band around 656 cm<sup>-1</sup> and is immersed in the background of noise.

## CONCLUSIONS

The appearance of the  $\nu_1$  mode in the  $B_g$  species contributing to  $\alpha_{xz}$ ,  $\alpha_{yz}$ , and  $\alpha_{xy}$  polarizability tensor components, the higher wavenumber values of the  $\nu_1$  mode, the observation of the inactive  $\nu_1$  and  $\nu_2$  modes in the infrared spectrum, and the lifting of degeneracies of the  $\nu_2$ ,  $\nu_3$ , and  $\nu_4$  modes suggest that the  $\text{SO}_4$  tetrahedron in LNHS is distorted confirming the X-ray data. The large splitting observed for the  $\nu_1$  mode in NHS and its decrease in NDS indicate that the  $\text{SO}_4$  tetrahedron is more distorted in NHS. Strong hydrogen bonds are formed between the H atoms of the  $\text{NH}_3$  and OH groups with the oxygen atoms of the  $\text{SO}_4$  group in both LNHS and NHS crystals. The presence of two bands for the nondegenerate symmetric bending mode ( $\nu_2$ ), the broadness of the stretching modes  $\nu_1$  and  $\nu_3$  of  $\text{NH}_3$ , and the appearance of two  $\text{NH}_3$  torsional bands at 90 K suggest the existence of two distinct  $\text{HONH}_3^+$  sites in the NHS crystal confirming the inelastic neutron scattering data. The two  $\text{NH}_3$  torsional modes observed at 90 K suggest that hydrogen bonds formed by  $\text{NH}_3$  protons become more pronounced at low temperature. In LNHS only one  $\text{NH}_3$  torsional mode is obtained and this indicates that  $\text{NH}_3\text{OH}^+$  sites in it are identical. ABC bands are observed in the OH stretching region in both the crystals, confirming the presence of strong hydrogen bonds.

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